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# **INVESTIGATIONS OF ULTRAVIOLET EMISSIONS** FROM THE IONOSPHERE

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PHILLIPS LABORATORY **Directorate of Geophysics** AIR FORCE MATERIEL COMMAND HANSCOM AIR FORCE BASE, MA 01731-3010



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This report describes the location, accessibility and utilization of the databases created from the measurements collected from the S3-4 satellite experiment, the Auroral Ionospheric Remote Sensor (AIRS) satellite experiment and the Horizon Ultraviolet Program (HUP) flown on Space Shuttles STS-4 and STS-39. Particular emphasis has been placed on comparing the predictions of the Atmospheric Ultraviolet Radiation Integrated Code (AURIC) with observed data.

The work of this contract involved the compilation and utilization of databases for satellite and Space Shuttle experiments conducted by the Phillips Laboratory/Geophysics Directorate/Ionospheric Modeling and Remote Sensing Branch (GPIM). The goal was to structure the results of these ionospheric measurements so that GPIM personnel and others could easily access these files to make ultraviolet background measurements and studies of the dynamics and composition of the auroral oval, the polar cap and the tropical belts. The auroral ovals are formed by the glowing ionosphere encircling each of the magnetic poles and are related to solar activity. The polar caps are not as bright as the ovals but have a number of features, such as sun-aligned arcs, that are the subject of current studies. The tropical belts are ionospheric emissions that circle the earth about 10 degrees poleward of the magnetic equator.

#### **ULTRAVIOLET BACKGROUNDS SATELLITE (S3-4)**

The raw telemetry tapes of the Ultraviolet Backgrounds experiment (S3-4) have been processed into vax-compatible tapes and are accessible at the Phillips Laboratory. Appendix 1 lists the files on these tapes for the revolution number (REV), day, format, and photometer filter. The S3-4 experiment made three simultaneous measurements. These were two spectrometer readings and one photometer measurement. The spectrometers scanned in the wavelength ranges 1070 A to 1930 A and 1620 A to 2900 A. The photometer had four bandpass filters that could be selected. These were centered at 1216 A, 1340 A, 1550 A, 1720 A and, also, no filter. The bandpass with no filter was that of the EMR 542G photomultiplier tube, 1070 A to 1930 A. From 3/22/78 to 9/10/78, data were collected from 355 full orbits and 447 partial orbits for a total of 728 hours of read-in time from a nominal altitude of 210 km and pointed in the nadir direction.

This database is being used to catalog the earth's ultraviolet glow and for scientific investigations of the ionosphere and auroral oval. GPIM and other laboratories believe there is still much information to be extracted from these data, which were collected in 1978. An example of the value of this database is the enclosed paper (Appendix 5) published in *Planet*. Space Sci., Vol. 40, No. 4, pp 481-493, 1992 by R. W. Eastes, R. E. Huffman and F. J. LeBlanc. This paper is a report on the nightglow emissions of atmospheric nitric oxide and molecular oxygen as seen in the ultraviolet from the S3-4 satellite.

#### HORIZON ULTRAVIOLET PROGRAM (HUP)

As part of the Horizon Ultraviolet Program, experiment 801A was flown on Space Shuttle STS-4. Also, a few hours of HUP data were collected on STS-39 with the instrument modified to collect data at 5 A wavelength resolution rather than 20 A as on STS-4. The HUP

instrument was designed to scan the earth's horizon at selected ultraviolet wavelengths, or to make spectral scans in the wavelength region between 1100 A and 1900 A. These data are stored on one side of an optical disc marked FPD002. A list of these data files can be found in Appendix 2, along with some of the fortran programs used in handling this data. The SETSDISC.DAT file, created from the STS-4 measurements, describes the disc arrangement of the data and lists the directories. The programs SETAVE.FOR, SETSVAL.FOR and GOSUN.FOR average, convert to radiance values, and assign these values to the tangent heights for the STS-4 data. These files were used for comparison with the AURIC predictions. Appendices 2.5 to 2.9 were written to manage the STS-39 data. SPECTRA.FOR extracts spectra from the files created by Boston College from the raw Lockheed tapes. HUPAVE.FOR averages the spectra. SENSCALC.PRO converts the counts to Rayleighs. HUPTIM.FOR writes the starting time of each spectra. SCAN40.FOR assigns radiance values for the HUP scan angles.

Because of personnel changes, Mr. LeBlanc became the principal point of contact for the HUP experiment being prepared to fly on Space Shuttle STS-39. He became involved in all phases of the project, including testing, integration, training, and flight operations. He was selected because a replacement was needed and he had performed similar functions in the past. These activities were very intensive from February 1990 until May 1991. Beyond this, the instrument sensitivities, wavelength, and scan-angle positions were redetermined. Results of this effort can be seen in the included paper "Horizon UV Program (HUP) Atmospheric Radiance Measurements" (Appendix 6).

#### **AURORAL IONOSPHERIC REMOTE SENSOR (AIRS)**

AIRS was designed to image the ionosphere between the east and west horizons at selected ultraviolet (UV) wavelengths from the PolarBEAR satellite and at times recorded the UV spectra looking in nadir. Data were collected primarily at northern latitudes, over the auroral oval and polar cap. The database is described in Appendix 3, along with the list of tapes containing AIRS spectra, the spectra files recorded on optical disc, and programs used to handle this data. The bulk of the image data have been placed on five double-sided optical disks. The unprocessed images can be accessed by keying in the desired observation times. This facilitates access to the data to such an extent as to make casual perusal feasible.

These databases (S3-4, HUP and AIRS) have been processed to obtain images, spectra, and background radiances of the earth's ionosphere and auroral oval. This was done for in-house studies, presentations, and papers, and also to answer requests from other laboratories. These occurrences have been documented in quarterly reports.

#### MODEL COMPARISON

The most recent use of these databases has been toward validation and modification of AURIC. AURIC predicts the atmospheric radiance for given look angles, altitudes, latitude and longitudes, 10.7-cm indices, universal times, and wavelengths. Appendix 4 contains information on the use of AURIC and associated programs. These files and programs are found in the

Auricsun workstation. The AURIC code is being used to predict horizon backgrounds for the oxygen 1356 A line (OI) and the Lyman Birge Hopfield (LBH) 1384 A and 1464 A bands. Figures 1 through 5 compare the results of AURIC, version 1.2, with HUP data. These horizon scans were selected to cover a range of solar zenith angles (SZA). The headings also identify the data set, the wavelength, and the scale factor used to make the AURIC model match the HUP data. Agreement with OI is reasonable, but the intensities of the LBH scans will need modification in the next version (AURIC 2.0).

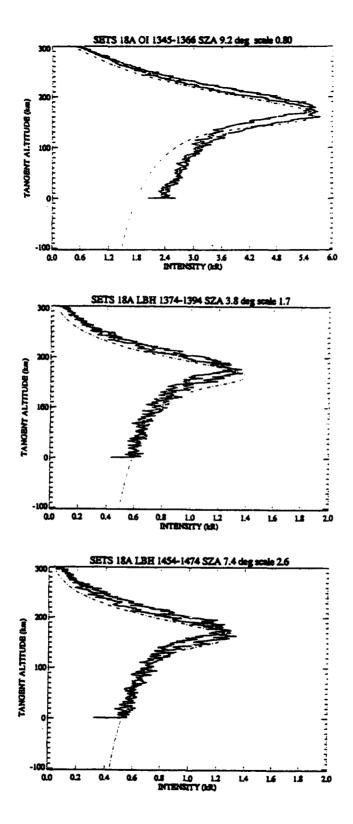


Figure 1. Upward and downward horizon intensity scans for the wavelengths indicated in the panel headings, observed from Space Shuttle STS-4 on Day 179, 1982 at 59400.0 UTsec and lat/lon = 24.5/280.0. The 10.7-cm solar index was 123.0. The dotted lines are the AURIC predictions modified by the scale factors in the headings.

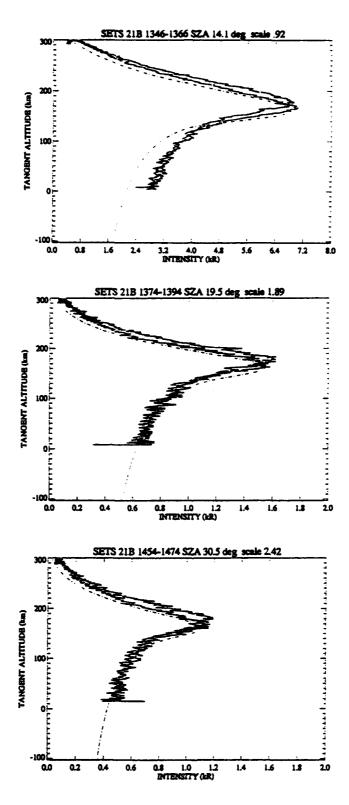


Figure 2. Upward and downward horizon intensity scans for the wavelengths indicated in the panel headings, observed from Space Shuttle STS-4 on Day 179, 1982 at 76918 UTsec and lat/lon = 16.8/233.7. The 10.7-cm solar index was 123.0. The dotted lines are the AURIC predictions modified by the scale factors in the headings.

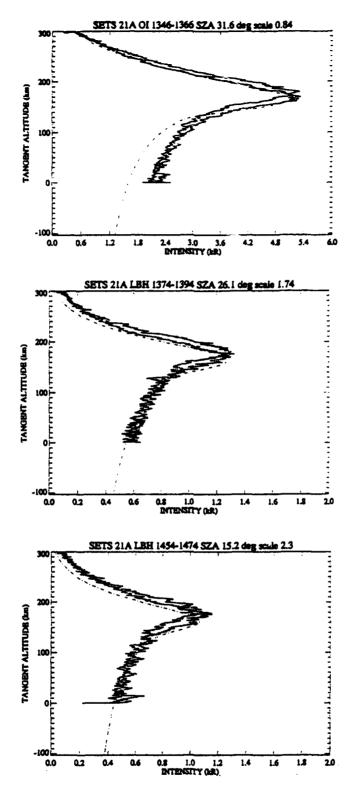


Figure 3. Upward and downward horizon intensity scans for the wavelengths indicated the panel headings, observed from Space Shuttle STS-4 on Day 179, 1982 at 76230 UTsec and lat/lon = 28.4/188.4. The 10.7-cm solar index was 123.0. The dotted lines are the AURIC predictions modified by the scale factors in the headings.

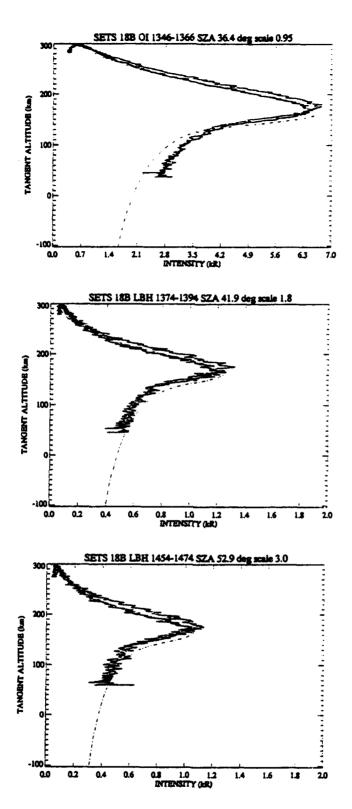


Figure 4. Upward and downward horizon intensity scans for the wavelengths indicated in the panel headings, observed from Space Shuttle STS-4 on Day 179, 1982 at 60970 UTsec and lat/lon = 7.6/321.0. The 10.7-cm solar index was 123.0. The dotted lines are the AURIC predictions modified by the scale factors in the headings.

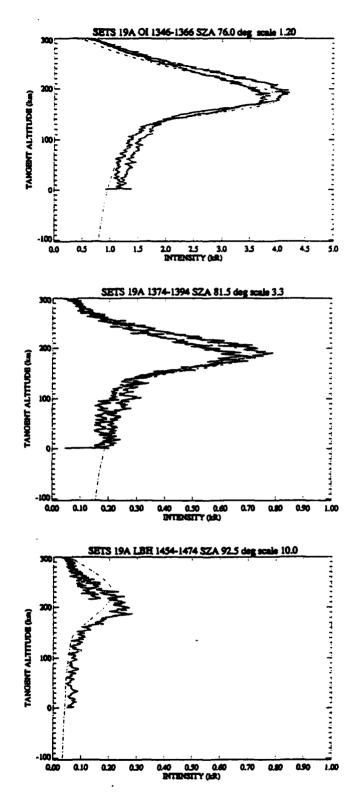


Figure 5. Upward and downward horizon intensity scans for the wavelengths indicated in the panel headings, observed from Space Shuttle STS-4 on Day 179, 1982 at 66955 UTsec and lat/lon = -10/328.0. The 10.7-cm solar index was 123.0. The dotted lines are the AURIC predictions modified by the scale factors in the headings.

APPENDIX 1: S3-4 Database

#### APPENDIX 1: S3-4 DATABASE

This table lists the raw telemetry tapes (SPECT. TAPE), the pre-processed spectrometer and photometer tapes (SPECT. P.P. TAPE and PHOT. P.P. TAPE) and the corresponding VAX compatible tapes. There are also columns for the orbit revolution number (REV), the month, day, year, telemetry format, and filter number. The number following the slash in tape identification numbers is the position of the file on the tape.

The data from the vehicle are Pulse-Code Modulated (PCM). Each mainframe consists of 120 words, 8 bits per word. The data, however, may be acquired at two data rates (32kbps and 64kbps). In general, the mainframe word locations for the designations are different for the two data rates. The 32kbps data rate is designated as Format A, and the 64kbps data rate is designated as Format C. A masterframe (one readout of each sub-commutated value) is 32 frames in Format A and 64 frames in Format C. The telemetry system may be summarized as follows:

8 bits/word
120 words/mainframe
960 bits/mainframe

Format	_A	Format C
32K	BITS/SEC	64K
33, 33	FRAMES/SEC	66.67
.03	SEC/FRAME	.015

\$3-4

#### SPECTROMETER AND PHOTOMETER

#### AGENCY, PRE-PROCESSED AND VAX COMPATIBLE TAPES

AGENCY S4S NT/800 BPI ,D-HD SPECTOMETER PRE-PROC. S4E NT/1600 BPI ,D-PE PHOTOMETER PRE-PROC. S4P NT/1600 BPI ,D-PE

VAX COMPATIBLE RM NT/6250 BPI ,D=GE (EXCEPT FOR (RM2885-RM2901), WHICH ARE NT/1600 BPI

1			s3-4				SPECT.	SPECT.	VAX COMP.	PHOT.	VAX COMP.
-		•	33-4				TAPE	P.P.TAPE	TAPE	P.P.TAPE	TAPE
1	REV	MTU	DAY	VD			IALE	I.I.IAID	IAL E	1.1.1.	
•	94	NO	DELL	IA	C	1	\$4\$343/1				
	99	3	22	78		3	S4S343/2	S4E136/1	RM2761/5	S4P555/1	RM2044/1
	164		26	78		3	545004/1	PARITY	ERRORS	UNABLE TO	PROCESS
	217	3	30	78		3	S4S343/5	S4E137/1	RM2762/5	S4P556/1	RM2045/1
	221	3	30		A	3	S4S342/1	S4E004/16	RM2913/5	S4P527/26	RM3513/6
	226		30	78	A	3	545022/1	S4E134/1	RM2895/3	S4P534/1	RM2784/1
	230	3	30	78	C	3	S4S138/1	S4E161/1	RM2625/13	S4P557/1	RM2046/1
	244	3	31	78		3	S4S118/1	S4E101/6	RM2904/2	S4P501/21	RM2784/6
	249	4	1	78		2	S4S333/1	S4E138/1	RM2759/6	S4P558/1	RM2047/1
	254	4	ī	78		4	S4S342/4	S4E036/11	RM2898/2	S4P570/41	RM2044/14
	258	4	ī	78		3	S4S333/3	S4E005/16	RM2896/1	S4P527/21	RM3513/5
	262	4	ī	-	A	2	S4S333/6	S4E118/1	RM2885/1	S4P530/41	RM3516/8
	266	4	2	78	C	3	S4S342/8	S4E139/1	RM2760/6	S4P559/1	RM2048/1
	270	4	2		A	3	S4S318/1	S4E001/6	RM2908/3	S4P526/11	RM3512/3
	274	4	2		A	3	S4S318/4	S4E002/6	RM2895/1	S4P527/11	RM3513/3
	277	4	_	78	A	3	S4S322/1		RM2901/3	S4P525/11	RM3511/3
	282	4	2	78	A	1	S4S322/4	S4E102/6	RM2759/7	S4P560/1	RM2049/1
	284	4	3 3		C	3	S4S322/5	S4E140/1 S4E141/1	RM2760/7	S4P561/1	RM2050/1
	287	4	3		A	3	S4S328/1	S4E019/11	RM2886/1	S4P530/16	RM3516/4
		4	3		A	3	S4S328/4	S4E019/11	RM2904/3	S4P531/16	RM3507/13
	291		3	78	A	3		· · · · ·		S4P562/1	RM2051/1
	295	_	_	78	C	_	S4S027/1	S4E134/6 S4E142/1	RM2895/4		RM2052/1
	301	4	4			2	S4S327/1		RM2761/6 RM2901/1	S4P563/1 S4P528/16	RM3514/4
	306	4	4		A	3	S4S327/3	S4E017/11		S4P528/16	RM3515/4
	308	4	4	78	A	3	S4S327/6	S4E018/11	RM2896/2	S4P529/16	RM3513/4
	312	4	4	78	A	_	545326/1	S4E015/11	RM2904/4	S4P527/16	
	318	4	5		C	3	S4S326/4	S4E143/1	RM2762/6	S4P564/1	RM2053/1 RM3511/4
	322	4	5	-	A	3	S4S326/7	S4E016/11	RM2904/5	S4P525/16	RM3513/2
	325	4	5		A	3	S4S301/1	S4E074/16	RM2904/6	S4P527/6	RM3514/2
	329	4	6		A	4	S4S301/4	S4E096/11	RM2667/3	S4P528/6	
	335	4	6	78	C	3	S4S301/8	S4E144/1	RM2761/7	S4P555/6	RM2044/2
	339	4	6	78	A	3	S4S321/1	S4E007/1	RM2903/1	S4P529/11	RM3515/3
	143	4	6		Ā	3	S4S321/4	S4E008/11	RM2897/1	S4P530/11	RM3516/3
	347	4	7	78	A	4	S4S321/7	S4E009/11	RM2898/1	S4P531/11	RM3515/10
	353	4	7	78	C	3	S4S117/1	S4E145/1	RM2762/7	S4P556/6	RM2045/2
	357	4	7		A	3	S4S332/1	S4E022/11	RM2887/1	S4P524/21	RM3510/5
3	360	4	7	78	A	3	S4S324/1	S4E012/11	RM2899/1	S4P523/16	RM3509/4

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                                 S4E117/1 RM2898/4
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APPENDIX 2: Reprint of Paper Entitled "NO and O<sub>2</sub> Ultraviolet Nightglow and Spacecraft Glow From the S3-4 Satellite"

# NO AND O2 ULTRAVIOLET NIGHTGLOW AND SPACECRAFT GLOW FROM THE S3-4 SATELLITE

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Abstract—Observations (1600-2950 Å) from the S3-4 satellite of the Earth's nightglow at  $\sim$ 6 Å resolution have been analyzed. These data indicate that the only significant emissions from NO are the  $\delta$  and  $\gamma$  bands. The relative brightness of these bands was consistent with previous observations at lower resolution. The data support the presence of O<sub>2</sub> emissions in the Herzberg I, II, and III bands in ratios ( $\sim$ 4.5:1:1.4), consistent with analyses of previous ground-based observations of these band systems. Spacecraft glow is seen from the N<sub>2</sub> Lyman-Birge-Hopfield bands, as previously reported; however, no spacecraft-related emissions were seen from NO or O<sub>2</sub>.

#### INTRODUCTION

Atomic oxygen is a minor constituent in the Earth's lower thermosphere; however, it is an important constituent due to the numerous emissions and chemical reactions in which it is involved. Unfortunately, it is difficult to measure the density of atomic oxygen in this region. Atmospheric drag limits the spatial and temporal coverage of in situ measurements. These limitations make remote sensing attractive; however. a thorough understanding of any emissions used is necessary since contributions from unresolved and unexpected emissions can lead to errors in interpreting low-resolution observations. One emission, which is proportional to the square of the atomic oxygen density, is from the O<sub>2</sub> Herzberg I bands (Dufay, 1941: Bates, 1954). The three-body recombination of atomic oxygen produces this emission, which is most prominent at night. The Herzberg I band emission peaks near 95 km (Packer, 1961; Stecher, 1965; Reed, 1968; L. Thomas et al., 1979; Llewellyn et al., 1979; Dickinson, 1980; Thomas and Young, 1981; Thomas, 1981; Greer et al., 1986; Murtagh et al., 1986).

Similar problems exist for measurement of atomic nitrogen (a minor constituent throughout the Earth's atmosphere). Radiative recombination of atomic nitrogen with atomic oxygen produces nitric oxide emissions which peak near 150 km (McCoy, 1983a, b). Few existing data sets have either simultaneous NO and O<sub>2</sub> emissions or sufficient sensitivity for high-resolution observations at night. However, data from the S3-4 satellite (Huffman et al., 1980) observed

emissions from both molecules at night with significantly higher resolution than most previous observations.

The nightglow spectra of NO and  $O_2$  observed from the S3-4 satellite and an analysis of them are presented in this paper. As an initial step, we demonstrate that these measurements are not compromised by the spacecraft glow seen on this satellite. Given sufficient understanding of the excitation mechanisms, the nightglow emissions from NO and  $O_2$  could be used for remote sensing of atomic oxygen and atomic nitrogen densities.

#### O2 emissions

Most analyses of the  $O_2$  Herzberg I band  $(A^3\Sigma_w^+ - X^3\Sigma_g^-)$  spectrum published (c.g. Slanger and Huestis, 1981) use ground-based observations (e.g. Broadfoot and Kendall (5.5 Å resolution) (1968); Chamberlain ( $\sim$ 0.8 Å) (1955, 1961); Krassovsky et al., ( $\sim$ 5 Å) (1962); Ingham (0.6 Å) (1962)). These observations cover the 3100-5000 Å region at  $\sim$ 1 to  $\sim$ 6 Å resolution; however, the brightest bands from the higher vibrational levels (v'=7-10) are below 3000 Å, wavelengths not observable from the ground.

Observation of thermospheric emissions at wavelengths shorter than 3000 Å requires that the instrument be on a satellite, sounding rocket or balloon. Such observations of O<sub>2</sub> emissions are scarce (Hennes, 1966; Huffman *et al.*, 1980; Cebula and Feldman. 1982, 1984; Sharp and Siskind, 1989). The data which have received the most analysis were from Hennes

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(1966), who observed the spectrum below 3000 Å at  $\sim 11$  Å resolution. Degen (1969) used these data in determining the relative vibrational populations of the A state.

After nearly half a century of controversy, the excitation of the  $O_2A^3\Sigma_n^+$  state is now attributed to three-body recombination of atomic oxygen (Bates, 1954; Barth, 1961, 1964; Thomas, 1981):

$$O(^{3}P) + O(^{3}P) + M \rightarrow O_{2}(A^{3}\Sigma_{n}^{+}) + M,$$
 (1)

which may be followed by

$$O_2(A^3\Sigma_n^+) \rightarrow O_2(X^3\Sigma_n^+) + hv$$

(Herzberg I bands). (2)

Quenching of the  $A^3\Sigma_n^+$  state could change either the total amount of emission or the vibrational distribution of the emission. Such changes could adversely affect the oxygen density derived from low-resolution observations (McDade et al., 1982).

As a test of theoretically predicted changes in the O<sub>2</sub> emission with altitude (McDade et al., 1982), a series of rocket observations, ETON (Greer et al., 1986; McDade et al., 1986a, b; Murtagh et al., 1986), designed to study the altitude dependence of the oxygen-dependent nightglow features, was launched. However, it was unable to conclusively determine the vibrational populations in these bands, or observe altitude variations in the Herzberg I vibrational populations. Higher-resolution observations may help in understanding these phenomena.

Emissions from  $O_2$  other than Herzberg I have been identified at wavelengths longer than 3000 Å. Using the Broadfoot and Kendall (1968) data. Slanger and Huestis (1981) showed that there is a contribution to the nightglow from the  $O_2c^3\Sigma_m$  state through Herzberg II band emissions; this system also emits below 3000 Å. More recently, Herzberg II band emissions and possibly emissions from the Herzberg III bands were identified in some ultraviolet data (2650–3050 Å) by Sharp and Siskind (1989).

#### NO emissions

Nitric oxide emissions observed in the Earth's atmosphere are generally attributed to either resonance fluorescence of sunlight or radiative recombination. At solar zenith angles and viewing geometries which exclude resonance fluorescence, radiative recombination is the only source of excitation, leading to emissions from .he v'=0 levels of the  $C^2\Pi$  and  $A^2\Sigma^+$  states which produce the NO  $\delta$  and  $\gamma$  bands respectively. Both band systems emit between 1600 and 2950 Å.

A number of daytime observations have been

reported since the first measurements by Barth (1964) (Barth, 1966; Pearce, 1969; Meira, 1971; Tisone, 1973; Tohmatsu and Iwagami, 1975, 1976; Thomas, 1978; Baker et al., 1977; McCoy, 1981, 1983a) with Cleary (1986) giving the most recent observations. These sounding rocket experiments measured emissions attributed to resonance fluorescence.

The number of observations at solar zenith angles large enough to prevent resonance fluoresence is more limited due to the weakness of the NO emission from radiative recombination; however, rocket-borne observations of this emission have been reported by Feldman and Takacs (1974), Sharp and Rusch (1981) and McCoy (1981, 1983a) near evening twilight. The number of observations further into night is smaller. Detection of NO  $\delta$  band emission in the nightglow was reported by Cohen-Sabban and Vuillemin (1973). Tennyson et al. (1986) reported more recent observations of NO emissions near local midnight.

The formation of the NO  $\gamma$  and  $\delta$  bands has been discussed in the literature (Sharp and Rusch, 1981; McCoy, 1983a, b). The radiative recombination excitation scheme for NO is best explained by the equations:

$$N(^4S) + O(^3P) \to NO(a^4\Pi) \to NO(C^2\Pi)(v'=0)$$
(3)

$$NO(C^2\Pi) \rightarrow NO(X^2\Pi) + hv (\delta \text{ bands})$$
 (4)

$$NO(C^2\Pi) \to NO(A^2\Sigma^+)(v'=0) + hv (1.22 \mu m)$$

(5)

$$NO(A^2\Sigma^+) \rightarrow NO(X^2\Pi) + hr$$
 (y bands). (6)

Since the a-C energy level crossing is below the NO dissociation level for only the v'=0 level of the C state, it is the only level energetically allowed. See McCoy (1981) for a detailed discussion of the allowed

#### MEASUREMENTS

The 53-4 satellite carried two 0.25 m Ebert-Fastie spectrometers which scanned the 1100-1900 Å and 1600-2950 Å regions simultaneously (Huffman et al., 1980) in  $\sim 22$  s. They were mounted on the three-axis stabilized spacecraft with their line-of-sight toward the nadir. Three pairs of slits were mounted on each spectrometer, enabling them to observe at  $\sim 1$ ,  $\sim 5$  or  $\sim 25$  Å resolution. Data were collected between March 30 and September 10, 1978. The elliptical, sun synchronous orbit had an inclination of  $96.5^{\circ}$  and took the satellite through the 160-270 km altitude region. The orbit crossed the Equator at approxi-

mately 10:30 and 22:30 local solar time. Apogee was in the southern hemisphere.

For these analyses, the data between 1600 and 2950 Å were summed. Although the nominal resolution of 5 Å full-width-half-maximum (FWHM) is confirmed by atomic emissions in the aurora, a coarser resolution was needed to fit the summed data. At least two effects contribute: first the resolution of the large number of scans summed will be degraded by any mechanical imperfections or wear over the time of the observations. The observations span  $\sim 3$  months. Second, the coarseness of the binning ( $\sim 1.3 \text{ Å}$ ) used for these data, relative to the 5 Å resolution, degrades the resolution in summed scans.

The wavelength calibration for the u.v. spectrometer was determined by using both the NO band emissions and the Rayleigh-scattered sunlight spectrum observed. The Rayleigh-scattered sunlight observations were compared with a high-resolution solar spectrum (Anderson and Hall, 1989) at wavelengths > 2300 Å. The resulting wavelength calibration agrees with the observed wavelengths to within  $\pm 0.75$  Å.

#### DISCUSSION OF DATA

Spacecraft glow considerations

Ultraviolet spacecraft glow emission from the  $N_2$  Lyman-Birge-Hopfield (LBH) bands was discovered on the S3-4 satellite (Huffman et al., 1980), following analysis of data from the u.v. spectrometers and the multiband photometer. When this emission was initially observed at night it was thought to be airglow. It was subsequently observed during the day, where it was much weaker than the normal dayglow. With the emission peaking at the v'=0 level of the upper state, it exhibited a different vibrational distribution than the normal LBH dayglow and auroral emissions.

This LBH band emission is now attributed to interaction of the spacecraft with the atmosphere (Conway et al., 1987). Using data from night-time measurements, they found that spacecraft glow brightness depended on altitude, and no emission was detected when the satellite was above 230 km. The altitude variation of the signal was proportional to the cube of the  $N_2$  density or the product of the square of the  $N_2$  density and the O density. Conway et al. (1987) concluded that the ambient nitrogen molecules have enough kinetic energy, relative to the spacecraft, to excite the lowest vibrational levels of the  $a^{\dagger}\Pi_{p}$  electronic state. Excitation mechanisms involving O or  $O_2$  were not excluded by the data, but no O I emissions were observed in the glow.

Additional suggestions regarding the production

mechanism of the spacecraft glow have come from Kofsky (1988). Swenson and Meyerott (1988), Cuthbertson and Langer (1989) and Meyerott and Swenson (1990). A heterogeneous reaction between two nitrogen atoms on a spacecraft surface is postulated, but the reactions producing the nitrogen atoms and many other details are in doubt. Controlled experiments rather than the current random observations would be valuable.

The data presented here are from the same time, but different orbits, as those analyzed by Conway et al. (1987). Therefore, the same spacecrast glow should be present. There is also the possibility of additional spacecraft glow at longer wavelengths where it would complicate analysis of the NO and O2 emissions. Previously, the longer-wavelength data from S3-4 have not been examined as thoroughly as the shorter-wavelength data for the presence of additional spacecraftrelated emissions. If there were enough energy available to produce the LBH band emissions, there would also be sufficient energy to produce NO and O2 emissions at longer wavelengths. Therefore, these data were first examined for additional sources of spacecraft glow by comparing two separate sums of the data from different altitude ranges.

Shown in Fig. 1 is a u.v. nightglow spectrum (1600–2950 Å) measured from the S3-4 satellite. This sum of 1447 individual spectra corresponds to observing each of the 932  $\sim$ 1.3 Å bins for 28.94 s. These data are smoothed over  $\approx$ 4 Å by taking running averages of 3 data bins. The  $\pm$ 1  $\sigma$  statistical error is shown for the NO  $\delta$  band (0, 2) peak at 2056 Å, and the NO  $\delta$  and  $\gamma$  band heads (v'=0) are marked. All the spectra summed were from solar zenith angles (SZAs) > 120°; had altitude > 230 km (average altitude = 251 km); and were outside the auroral regions.

Spectral scans containing auroral emissions were excluded by the following criteria. Whenever enhanced atomic oxygen emissions (1304 and 1356 Å) were observed by the far-u.v. spectrometer, which simultaneously observed shorter wavelengths, the data were excluded from the sums shown in this paper. These oxygen emissions are readily apparent whenever significant auroral emission occurs. The selection criteria is considered sufficient since N<sub>2</sub> Vegard–Kaplan and LBH band emissions, which are also produced in the aurora, are absent from the data shown in Fig. 1.

A second sum of 633 spectra observed when the satellite was at lower altitudes (<220 km (average altitude = 210 km)) is shown in Fig. 2. These observations, like those above, were from non-auroral conditions but with SZAs > 110°. The  $\pm 1~\sigma$  statistical error is shown for the emission peak at 2056 Å. The

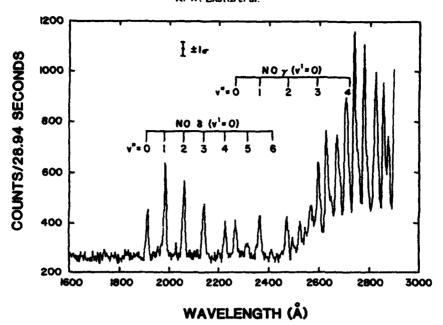


Fig. 1. Ultraviolet nightglow spectrum measured from S3-4 at the highest altitudes. This sum of 1447 individual spectra, each taking  $\sim 22$  s, corresponds to observing each of the 932  $\sim 1.3$  Å bin for 28.94 s. All spectra summed were from solar zenith angles  $> 120^\circ$ : altitudes > 230 km (average altitude = 251 km); and outside the auroral regions. These data are smoothed over  $\approx 4$  Å by taking running averages over the 3 data bins. The  $\pm 1$   $\sigma$  statistical error is shown for the NO  $\delta$  band (0, 2) peak at 2056 Å, and the NO  $\delta$  and  $\gamma$  band heads ( $\sigma$  = 0) are marked.

dashed curve represents the response of the spectrometer to a 2 R emission line as a function of wavelength (after subtracting 100 counts per bin of background signal).

Weak emissions from the LBH spacecraft glow were seen at the shorter wavelengths of the loweraltitude spectrum (Fig. 2). This emission brightness is consistent with the N<sub>2</sub> LBH band emissions from spacecraft glow analyzed by Conway et al. (1987). Conway found that for an altitude of 210 km, 20-50 R were seen in the 1400-1700 A region. Assuming the vibrational distribution derived by Conway et al. (1987), the data in Fig. 2 at 1650 Å indicate that the LBH band emissions in the 1400-1700 Å region are ≤ 160 R. Uncertainties in the background make a more precise brightness difficult to obtain. Figure 4 shows both the summed low-altitude data and the LBH spectrum which the instrument would record, assuming the vibrational population derived by Conway et al. (1987) (v' = 0 to 6 populations of 1.0, 0.76. 0.18, 0.16, 0.39, 0.20 and 0.00, respectively). As shown by Fig. 2, emissions from the LBH bands should be insignificant at the wavelengths of the NO  $\delta$  and NO y band emissions in Figs 1-3.

In Figs 1 and 2 the NO  $\delta$  (v' = 0), NO  $\gamma$  (v' = 0),

and  $O_2$  Herzberg I (v'=4-11) bands are conspicuous. The presence of other  $O_2$  bands will be discussed later. The NO bands dominate the spectrum at wavelengths <2500 Å, while the  $O_2$  Herzberg bands dominate at longer wavelengths. Within statistical error the brightness of the  $O_2$  Herzberg bands was the same in both sums. This was expected. The three-body recombination of  $O_2$ , which produces the  $O_2$  Herzberg band emission, peaks near 95 km and all significant emission comes from below 160 km, the typical perigee of the satellite. This indicates that no statistically significant emission comes from spacecraft glow at wavelengths 2500 Å <  $\lambda$  < 2950 Å.

On the other hand, the brightness of the NO emission changes. A comparison of the higher altitude (Fig. 1) and lower altitude (Fig. 2) data shows that the NO emissions observed from higher altitudes are 1.6 times brighter than at lower altitudes. However, the NO emissions are expected to peak at  $\sim 150$  km. Since the lower altitude data in Fig. 2 come from altitudes of 160-220 km, more of the emission is above the satellite and not in the field-of-view of the nadirviewing instrument. The brightness of the NO bands decreases approximately as expected according to the model of McCoy (1983b). This agreement indicates

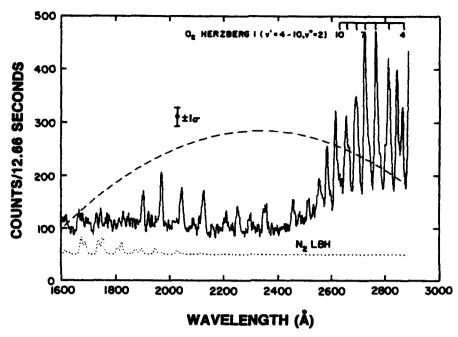


Fig. 2. Ultraviolet nightglow spectrum measured from S3-4 at the lowest altitudes. Shown here is the sum of 633 spectra observed from lower altitudes (<220 km (average altitude = 210 km)) than the data in Fig. 1. These observations are from non-auroral conditions with solar zenith angles > 110". These data are smoothed over  $\approx$ 4 Å by taking running averages over the 3 data bins. The  $\pm$ 1  $\sigma$  statistical error is shown for the emission peak at 2056 Å. The smooth, dashed curve represents the response of the spectrometer to a 2 R emission line as a function of wavelength (after subtracting 100 counts per bin of background signal). The dotted curve shows the N<sub>2</sub> Lyman-Birge-Hopfield band emissions calculated using the relative vibrational populations from Conway et al. (1987) (v'=0 to 6 populations of 1.0, 0.76, 0.18, 0.16, 0.20, 0.00, respectively) and scaled to match the data near 1650 Å.

that there were probably no NO emissions produced by spacecrast glow. A more definitive conclusion is not justified from this comparison since the observations cover about 3 months, during which the NO production rate (and emission) may vary. However, spacecrast glow should not interfere with the analysis of the NO emissions.

Comparison of the data in Figs 1 and 2 indicates that there are no significant spacecraft-related emissions from NO or  $O_2$ . Therefore, the two sums were combined to improve the counting statistics. Shown in Fig. 3 is the sum of all scans shown in Figs 1 and 2 plus scans at altitudes between 220 and 230 km. The resulting sum of 2385 scans will be used for the remaining analysis. The  $\pm 1 \sigma$  statistical error is shown for the emission peak at 2056 Å.

#### Analysis of NO band emissions

Before spectral analysis the background signal must be subtracted. To accomplish this the background between the NO  $\delta$  and  $\gamma$  bands was first fit to a straight line. The emission observed from the N<sub>2</sub> LBH bands,

while insignificant, was avoided. At wavelengths where overlapping emissions from O<sub>2</sub> prohibited direct observation of the background, a constant background was used. This constant background level was determined during the fitting of the molecular emissions. The constant level was consistent with the fit to shorter wavelengths. In Fig. 4 the data shown in Fig. 3 has been plotted as brightness versus wavelength after subtraction of the background signal.

First, a synthetic NO  $\delta$  band spectrum was fitted to the data (1877-2429 Å) in Fig. 4. The synthetic spectrum calculation for NO used Franck-Condon factors from Nicholls (1964) for the  $\delta$  bands and from Spindler et al. (1970) for the  $\gamma$  bands. Molecular constants from Huber and Herzberg (1979) were used for the C and X states, and Engleman and Rouse (1971) were the source of constants for the A state. The synthetic spectrum calculations used only the  $P_1$ ,  $Q_1$ , and  $R_1$  branches. These are sufficient at  $\sim 5$  Å resolution, given the noise in the data and the coarseness of the binning. The band origins calculated for both the  $\delta$  and  $\gamma$  bands were in good agreement ( $\pm 0.2$  Å)

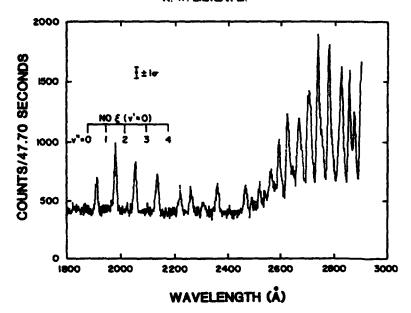


Fig. 3. Ultraviolet Nightglow spectrum measured from S3-4 at all altitudes. This spectrum is the sum of all spectra shown in Figs 1 and 2 plus spectra at altitudes between 220 and 230 km. The resulting sum contains 2385 scans. These data have not been smoothed as were the data in Figs 1 and 2. The  $\pm 1~\sigma$  statistical error is shown for the emission peak at 2056 Å.

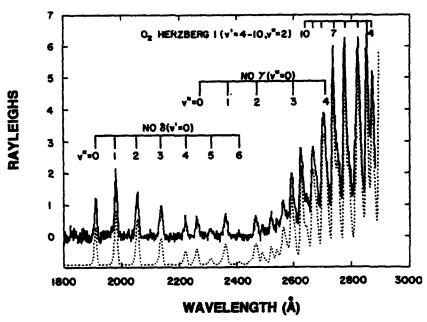


Fig. 4. Ultraviolet nightglow spectrum in Rayleighs as measured from S3-4. After subtraction of the background from the data shown in Fig. 3, the relative sensitivity of the instrument was removed. The NO  $\delta$  and  $\gamma$  bands are marked. The dotted curve is the best fit to the data when using the NO  $\delta$ , NO  $\gamma$ , O<sub>2</sub> Herzberg I, and O<sub>2</sub> Herzberg II bands.

with published values by Wallace (1962) or Pearse and Gaydon (1963).

The synthetic spectrum included the rotational structure evident in the observed bands. The spectral resolution and rotational temperature  $(T_r)$  were varied iteratively to determine the least-squares fit. The singular value decomposition method (e.g. Press et al., 1986) was used to solve for the least squares in all fitting to the data from the S3-4 satellite. The best fit (shown in Fig. 5) was obtained using 5.9 Å resolution (FWHM) and a rotational temperature  $(T_r)$  of  $\sim 700$  K, approximately the ambient temperature (Hedin, 1987) at the altitude of the NO emission peak (McCoy, 1983b) of  $\sim 150$  km.

Next, the NO  $\gamma$  bands were fitted using data (2234–2478 Å) from the (0, 0), (0, 1), and (0, 2) bands. Just as for the  $\delta$  bands, the spectral resolution and rotational temperature of the  $\gamma$  bands were iterated until a best fit was obtained. A resolution of  $\sim 5.9$  Å and  $T_r$ , of  $\sim 700$  K (same as for the  $\delta$  bands) fitted the data best. The best fit is shown in Fig. 5.

The NO  $\gamma$ : NO  $\delta$  brightness ratio of 1:1.4 (based on the best fit and Franck-Condon factors for the (0, 1) bands) indicates that the branching ratio from the  $C^2\Pi$  to the  $A^2\Sigma$  state is  $0.41\pm0.16$  (see equations (4) (7)). Within experimental error this agrees with previous observations (0.30±0.06 by Tennyson et al., 1986; 0.21±0.06 by Sharp and Rusch, 1981; 0.23±0.05 by McCoy, 1983a).

The emission from the NO  $\varepsilon$  bands was <0.25 R ( $\Sigma_r$ -I(v'=0, v'')). The NO  $\varepsilon$  (v'=0) band heads are shown in Fig. 3. This limit would be lower but for the LBH glow from the spacecraft.

#### Analysis of O2 band emissions

After fitting the NO emissions, the longer-wavelength data were fit to synthetic molecular spectra of the  $O_2$  Herzberg I  $(A^3\Sigma_n^+ - X^3\Sigma_g^-)$  and  $O_2$  Herzberg II  $(c^1\Sigma_n^- - X^3\Sigma_g^-)$  bands. Possible contributions by the  $O_2$  Herzberg III  $(A'^3\Delta_n - X^3\Sigma_g^-)$  bands will be discussed later since they have never been observed conclusively in the nightglow. The relative vibrational population of each level of the A state was fitted to the data independently, but the relative vibrational populations of the c state were fixed. It was necessary to assume relative populations for the c state due to the large number of parameters being fitted, resolution of the spectrum, statistical uncertainty, and the brightness of the Herzberg II bands relative to the Herzberg I bands.

To generate the synthetic  $O_2$  Herzberg spectra, molecular constants suggested by Krupenie (1972) for the X state were used with constants from Ramsay (1986), Borrell et al. (1986), and Coquart and Ramsay (1986) for the c, A, and A' states. Transition probabilities from Bates (1989) were used for the Herzberg I, Herzberg II, and Herzberg III bands. Rotational formulae from Degen (1969) (for the Herzberg I

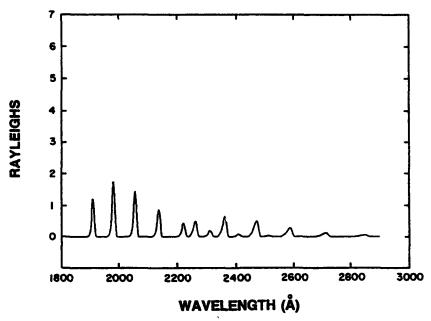


Fig. 5. Best fit to the NO  $\delta$  and  $\gamma$  bands in the summed nightglow data shown in Fig. 4. Shown in this plot is the best fit of the NO bands to the summed nightglow data shown in Fig. 4.

bands) and Kovacs (1969) (for Herzberg II and Herzberg III bands) were used. The rotational temperature of 200 K was based on the temperature of the MSIS neutral atmosphere (Hedin, 1987) at  $\sim$ 95 km. The resolution should be consistent with the shorter wavelength NO emissions; therefore, the  $O_2$  spectra were calculated at 5.9 Å resolution.

The fitting was done iteratively using the last vibrational populations for the A state to determine those for the c state until the populations ceased to vary. The relative vibrational populations for the c state were derived from the populations of the A state by the relationship:

population(
$$v'_{\epsilon}$$
) = population( $v'_{A} + 2$ ),

an approach used by Slanger and Huestis (1981). Since vibrational populations for  $v' \leq 3$  of the A state were not derived from this data, the populations for  $v' \leq 5$  of the c state were taken from Slanger and Huestis (1981). The resulting fit, using the NO  $\delta$ , NO  $\gamma$ , O<sub>2</sub> Herzberg I, and O<sub>2</sub> Herzberg II bands, is shown in Figs 4 and 6. In Fig. 4 the synthetic spectrum is plotted as a dashed line beneath the data. Shown in Fig. 5 is the contribution of the NO  $\delta$  and  $\gamma$  bands. The total fit and the individual contributions from the O<sub>2</sub> Herzberg I and Herzberg II bands are shown in Fig. 6. The Herzberg I is plotted as a dashed line,

offset beneath the total fit. The Herzberg II spectrum is shown by the solid line offset beneath the total fit.

Although they were not clearly resolved, the Herzberg II bands have previously been observed in the nightglow. Using the Broadfoot and Kendall (1968) data Slanger and Huestis (1981) showed there is a contribution to the nightglow from the  $O_2c^1\Sigma_r^-$  state through Herzberg II band emissions. They found a Herzberg 1: Herzberg 11 brightness ratio of 4:1 (when integrated over all emission, including wavelengths longer than 3000 Å), and showed that the Herzberg II bands appear as shoulders on many of the Herzberg I bands in the 3000-4000 Å range. More recently Sharp and Siskind (1989) reported observations at wavelengths < 3050 Å where the brightness ratio was 9:1. The Herzherg 1: Herzherg II brightness ratio from the observations shown here is  $3.7\pm0.1:1.0\pm0.3$  for Herzberg 1 emissions from the v'=4 to 10 levels, over the entire range of emission. This ratio would be larger if the lower vibrational level emissions (v' < 4) from the Herzberg I system were included. However, the populations for levels from v' < 4 were uncertain, or not determinable. Assuming that the vibrational levels with v' < 4 were populated, relative to v' = 4, according to the populations found by Degen (1989) or Slanger and Huestis (1981), the brightness ratio would be ~4.8:1.

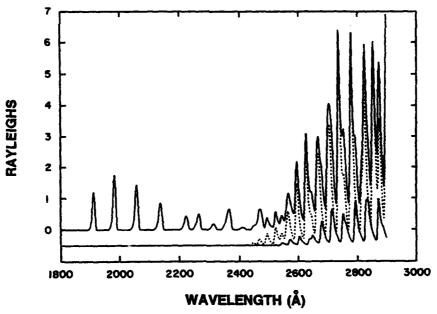


Fig. 6. Best fit to the nightglow data shown in Fig. 4 (upper solid histogram). This fit uses the NO  $\delta$  and  $\gamma$  bands as well as the O<sub>2</sub> Herzberg I (v'=4 to 10 only) and II bands. The dotted line shows the contribution of the Herzberg I bands and the bottom solid line shows the contribution of the Herzberg II bands. For this spectrum  $T_r$  is 200 K and the resolution is 5.9 Å.

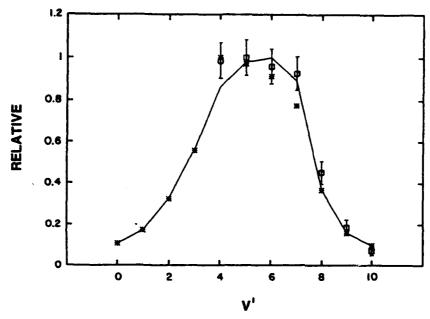


Fig. 7. Relative vibrational level populations for the  $O_2$   $A^3\Sigma_0^+$  state, as derived from the Herzberg I emissions, are plotted as squares with error bars. The populations derived by Degen (1969) are plotted as \*. and the populations derived by Slanger and Huestis (1981) are represented by the solid line.

Shown in Fig. 7 as squares with  $\pm 1$   $\sigma$  error bars are the relative vibrational populations for the  $O_2A^3\Sigma_n^+$  state as determined from the Herzberg I emissions. The populations given for these data are based on the assumption of equal emission probability from all levels. For equal populations the total emission from each level would be equal, under this assumption.

Vibrational levels with  $v' \leq 3$  have such a small contribution to the wavelength range scanned that they could not be determined. Emissions from the Herzberg I v' = 11 level could not be fit since the emission has no distinct peaks at the wavelengths scanned. Populations for the  $A^3\Sigma_n^+$  derived by Slanger and Huestis (1981) and Degen (1969) from analysis of the data of Broadfoot and Kendall (1968) and Hennes (1966) are shown for comparison as a solid line and asterisks, respectively. When fitting the Herzberg I and Herzberg II emissions in the data, the Herzberg I relative vibrational populations which gave the best fit were  $0.99\pm0.08$ ,  $1.00\pm0.08$ ,  $0.96\pm0.08$ ,  $0.93\pm0.08$ ,  $0.45\pm0.06$ ,  $0.19\pm0.04$ , and  $0.07\pm0.02$  for v'=4 to 10, respectively.

The populations for the v'=4 to 9 vibrational levels agree well with previous observations. The v'=10 contribution found may be subject to sig-

nificant systematic errors. This level was sensitive to the molecular constants. It was also more sensitive to the background subtraction than the other levels. The relative population of  $0.07\pm0.02$  shown in Fig. 7 approximately doubles if the wavelength-dependent background subtraction used for the NO emission region is extrapolated to longer wavelengths. This suggests that the v'=10 population derived here is best considered as a lower limit. The contributions from individual vibrational levels (v'=4 to v'=10) is shown in Fig. 8.

The recent publication of transition probabilities for the  $O_2$  Herzberg III band emission (A'-X') (Bates, 1989) aids calculation of the spectrum of this band system. This band system has never been conclusively observed in the upper atmosphere. Sharp and Siskind (1989) have reported the possible presence of it in some observations at wavelengths < 3050 Å. Observation of this band system is difficult since the wavelengths of emission closely match those of the Herzberg I system from the A state and the emissions are weak.

Inclusion of Herzberg III emissions in the fitting of these data does not significantly affect the relative populations derived for the A state (relative populations are  $0.95\pm0.08$ ,  $1.00\pm0.8$ ,  $0.87\pm0.08$ ,

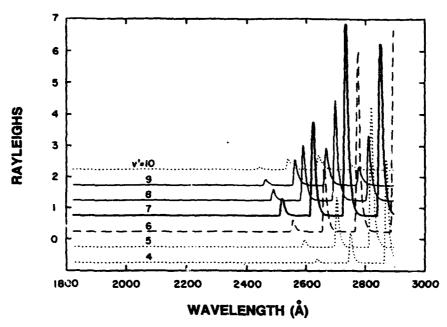


Fig. 8. Best fit of the v'=4 to  $10~O_2$  Herzberg I bands to the data in Fig. 4 when using only the  $O_2$  Herzberg I and II bands in the fitting (in addition to the NO bands). Shown here are the contributions from each v' to the total shown in Fig. 6.

 $0.86 \pm 0.08$ ,  $0.38 \pm 0.05$ ,  $0.20 \pm 0.04$ , and  $0.06 \pm 0.02$ for v'=4 to 10, respectively). It does slightly decrease the derived brightness (by  $\sim 0.9$ ). The goodness of fit improves significantly also. The reduced  $\chi^2$  decreased from 5.8 to 4.7 when fitting to the data in bins 600-958 (2419-2896 Å). The populations for the  $c'\Sigma$  state (Herzberg II) were handled in the same way as when fitting Herzberg I and Herzberg II without Herzberg III. The relative populations used for the  $\Lambda^{\prime 3}\Delta$  state (Herzberg III) were from Slanger and Huestis (1981). The fit to the data is shown in Fig. 9. The Herzberg I (dotted), Herzberg II (solid), and Herzberg III (dashed) components of the fit are shown separately as the offset plots beneath the total fit. The brightness ratio for Herzberg I: Herzberg II: Herzberg III emissions in these data was  $3.4\pm0.1:1.0\pm0.1:1.4\pm0.5$ for Herzberg I v'=4 to 10. Assuming that the vibrational levels with v' < 4 are populated, relative to v' = 4, according to the populations found by Degen (1968) or Slanger and Huestis (1981), the brightness ratio would be ~4.5:1:1.4, which agrees with previous observations at longer wavelengths but is significantly lower than other observations (Sharp and Siskind, 1989) in this wavelength range (<2950 Å).

The identification of Herzberg III emissions in these data must also be considered tentative as was Sharp and Siskind's (1989) identification. Any errors in cal-

culating the Herzberg I band emission, which is blended with the Herzberg III emission, could significantly affect the conclusions. Errors are possible due to the accuracy of the Dunham coefficients used and the approximations assumed when using them in calculating the spectra (see Slanger and Cosby, 1988). A significant portion of the emission in this wavelength region involves higher vibrational levels for which the molecular constants are less well known than for the lower levels.

#### CONCLUSIONS

The nightglow in the 1600-2950 Å region has been analyzed to aid in developing remote-sensing methods for atomic oxygen and atomic nitrogen. The NO  $\delta$  and  $\gamma$  band emissions from radiative recombination have been observed at sufficient resolution (5.9 Å) to derive a rotational temperature of  $\sim 700$  K. The  $C^2\Pi$  to  $A^2\Sigma$  branching ratio observed was  $0.41 \pm 0.16$ , which is in agreement with previous observations. Emission from the NO z bands was < 0.25 R.

Vibrational populations for the  $O_2$  A state (0.95, 1.00, 0.87, 0.86, 0.38, 0.20, 0.06 for v'=4 to 10, respectively) derived from these Herzberg I band observations are in agreement with previous observations. Emissions from the Herzberg II and III bands were

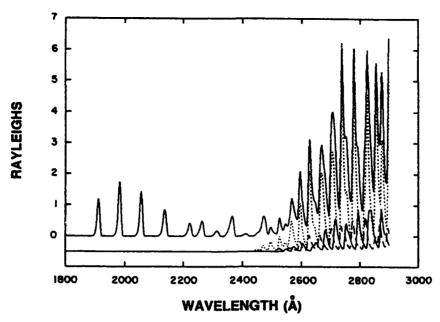


FIG. 9. THE UPPER RISTINGRAM SHOWS THE SUMMED NO AND O2 emissions (Herzberg I, II, and III) which best fit the data in Fig. 4.

The Herzberg I, II, and FI emissions are shown separately below the total as dotted, solid and dashed lines, respectively.

also suggested by the data. The brightness ratio for Herzberg I: Herzberg II: Herzberg III emissions in these data was  $\sim 4.5:1:1.4$ , which agrees with previous observations at longer wavelengths but is significantly lower than other observations in this wavelength range (<2950 Å).

No spacecrast-related emissions were detectable from the NO or  $O_2$  band emissions observed. Weak emissions from the  $N_2$  LBH bands were seen. These are attributed to spacecrast interactions, as previously reported. The lack of spacecrast-related emissions from  $O_2$  and NO may constrain the excitation mechanisms for LBH emission from spacecrast.

Acknowledgements—We would like to thank Dr J. Thomas for enlightening discussions on calculations involving the NO y bands and Dr J. Retterer for his suggestions on more robust fitting methods. We are also indebted to the reviewer for helpful and clarifying suggestions.

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APPENDIX 3: Reprint of Paper Entitled "Horizon UV Program (HUP)
Atmospheric Radiance Measurements," Describing HUP
Measurements

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# Horizon UV Program (HUP) Atmospheric Radiance Measurements

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#### **ABSTRACT**

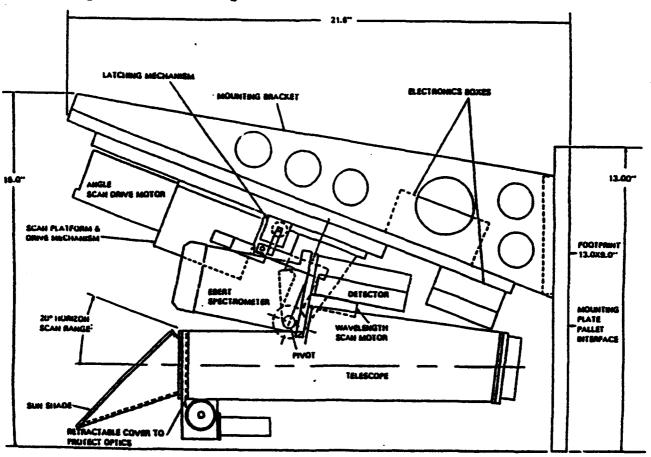
Spectrometric observations (1100 Å to 1800 Å) of the ultraviolet radiance of Earth's atmosphere have been obtained by the Horizon Ultraviolet Program (HUP) sensor from Space Shuttle altitude at two different points in the solar cycle. Most of the radiance measurements were made during horizon scans of the earthlimb or while the HUP field of view was directed toward nadir. The expected variation of dayglow radiance with solar zenith angle (SZA) in the several spectral bands was observed. These results compare favorably to the predictions of a preliminary version of the Atmospheric Ultraviolet Radiance Integrated Code (AURIC).

### 2. INTRODUCTION

A small but growing body of experimental measurements of atmospheric ultraviolet radiance is now extensive enough that it is reasonable to compare those data, obtained under a variety of geophysical conditions, with the predictions of atmospheric radiance models. This paper presents some results from just such a comparison. The range of conditions for which comparison is being reported here is purposely limited to simple cases of ultraviolet dayglow; at this level, issues of radiation transport and optically thick emission features are being avoided. We concentrate on the optically thin molecular nitrogen Lyman-Birge-Hopfield bands and the atomic oxygen line at 1356 A. These provide an adequate set of initial test cases. In addition it is precisely those features of the dayglow which are candidates for use in the remote determination of the distribution of electrons and atmosphereic species in the ionosphere. Spatial and temporal variability are important background characteristics for any surveillance systems but they are not addressed in this paper. There are other potential byproducts such as monitoring solar activity for its effect on communications and radar and as a horizon sensor to monitor spacecraft stability.

#### 3. EXPERIMENTAL

Figure 1 is a diagram of the HUP instrument. It was an eighth meter Ebert-Fastie spectrometer fitted with a 230mm telescope. This was mounted on a platform which allowed the instrument to be pivoted through 20 degrees in 41 seconds. It could be operated as a photometer at a selected wavelength or as a spectrometer recording a spectrum every 5.8 seconds. The instrument was mounted in the shuttle bay so that it could be made to point perpendicularly out of the bay or scan 20 degrees between vertical and towards the shuttle nose so that when shuttle was in its gravity gradient mode (nose down) the horizon would be scanned. The diagram shows HUP oriented as it would be to view vertically out of the bay.



HORIZON ULTRAVIOLET PROGRAM

**FLIGHT SENSOR** 

Fig. 1. Schematic of HUP instrument shown as mounted on Shuttle in launch position. The telescope, spectrometer and detector can be pivoted as a unit.

Fig. 2. shows the measured calibration curves for both flights for the wavelength range of 1150 A to 1850 A, shown in counts per Rayleigh-sec. The difference is due to the slit change, with a smaller contribution to loss of sensitivity after nine years of aging.

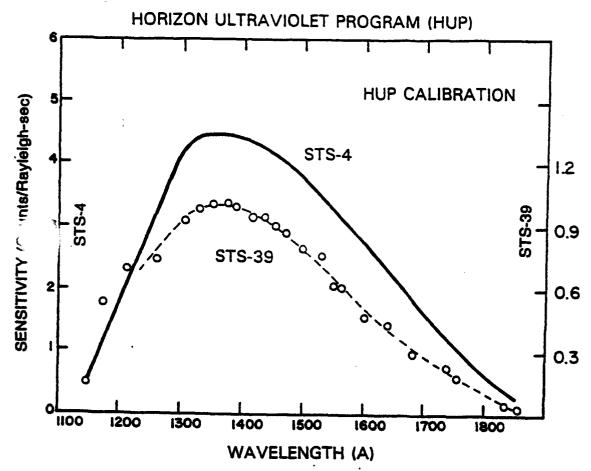


Fig. 2. HUP instrument sensitivity calibration for Shuttle flights.

For STS-39, the second flight, the slits, Table 1, of the spectrometer were narrowed to increase the spectral resolution. This also narrowed the FOV. In addition, the instrument was programmed to concentrate on oxygen 1356 Å and LBH radiation at 1384 Å to provide suitable data for comparison with ionosphereic models such as AURIC.

The measured wavelength resolution for STS-4 was 20 Å and 8.3 Å for STS-39. STS-4 had a 0.2  $\times$  1.5 deg field of view (FOV) with a footprint 3.8  $\times$  28 km at a distance of 1100 km. For STS-39 the the FOV was 0.06  $\times$  1.5 deg with a footprint of 1.2  $\times$  28 km at 1100 km. STS-4 flew at 300 km altitude between latitudes of 28.6 degrees and STS-39 flew at 259 km between 57 degrees latitude.

#### 4. Results

The observations presented here result from data collected on STS-4 June 28 and 29 1982 when the 10.7 cm solar index was about 112 and on May 5 1991 (STS-39) when the 10.7 cm solar index was about 183.

Fig. 3. displays a spectrum seen looking down from STS-39. The atomic lines are clearly resolved while the Lyman-Birge-Hopfield band system of Nitrogen (LBH) covers an extended wavelength range at lower

intensities. Fig. 4 is a spectrum from STS-4. The twenty Angstrom resolution is seen in the broader atomic line spectra.

Table 1. Features of the HUP instrument as flown.

	STS-4	STS-39
Slit Width:	0.80 mm	0.24 mm
Resolution:	20 Å	8.3 Å
Field of View:	0.188 x 1.45 deg.	0.058 x 1.45 deg
Footprint; 1100 km:	3.8 x 28 km	1.2 x 28 km
Wavelengths:	1216 Å H 1304 Å O 1356 Å O, N <sub>2</sub> LBH 1384 Å N <sub>2</sub> LBH 1414 Å N <sub>2</sub> LBH, N 1464 Å N <sub>2</sub> LBH 1493 Å N <sub>2</sub> LBH, N 1554 Å N <sub>2</sub> LBH	1356 Å O 1384 Å N <sub>2</sub> LBH
Maximum latitude:	28.6 deg	57 deg
Altitude:	300 km	268 km

The AURIC model, described by Strickland, <u>et al</u> else-where in this volume has been used to compare with our observations. We have overlaid our data with the AURIC predictions for the appropriate day, time of day, longitude, latitude, planetary index  $A_p$  and 10.7 cm solar activity index.

Figs. 5. to 9. display horizon scans of oxygen 1356 Å radiation obtained from STS-4 for a range of solar zenith angles with AURIC predictions superimposed (dotted line). The HUP traces are double, showing the upward and downward horizon scans. The offset seen between the upward and downward scans could be caused by a small amount of backlash in the cam groove driving the scan platform or by an offset of the zero angle position for the barrel cam rotation. Since both traces lie within the accuracy of our knowledge of shuttle pointing we have deferred fine tuning thes results until it becomes advantageous to do so. With the exception of the night measurement, the agreement between tangent height (about 165 km) of the maximum radiance as well as the magnitude of the peak brightness (5 to 6kR) with the prediction of the model is very gratifying. It provides reassurance that, as expected, the model is fundamentally sound and that HUP, Polar BEAR and other good experimental data will be invaluable in validating AURIC for more problematic conditions.

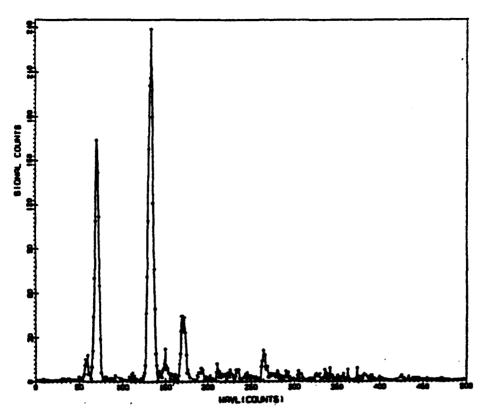


Fig. 3. HUP daytime spectra looking down from Space Shuttle STS-39 with 8.3A resolution.

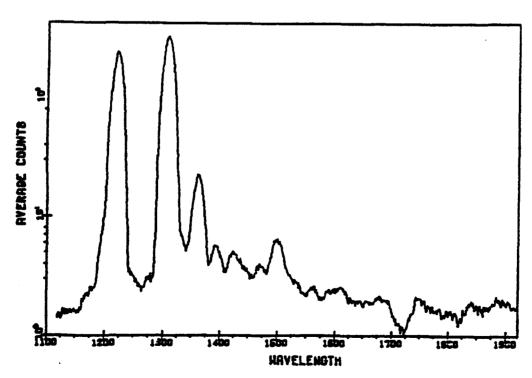
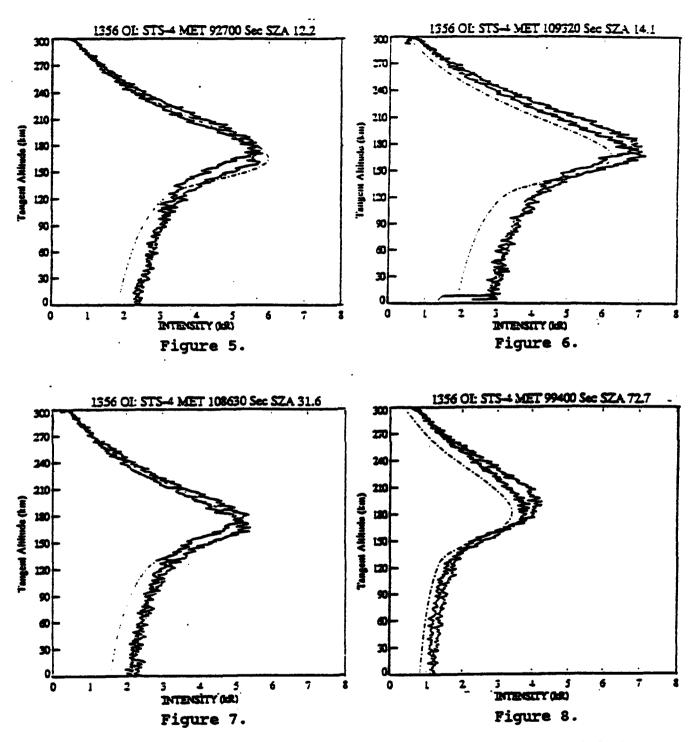


Fig. 4. HUP daytime spectra looking down from Space Shuttle STS-4 with 20A resolution.



Figs. 5 to 8. HUP horizon scans (up and downward) of 1356 A Oxygen emission observed from Space Shuttle STS-4 for a range of solar zenith angles, overlaid with the AURIC prediction (broken line).

Fig. 10 is a plot of the peak 1356 Å oxygen line intensity as seen from STS-4 in the nadir direction. This was done for a range of SZAs (crosses) and is overlaid with AURIC predictions. The differences

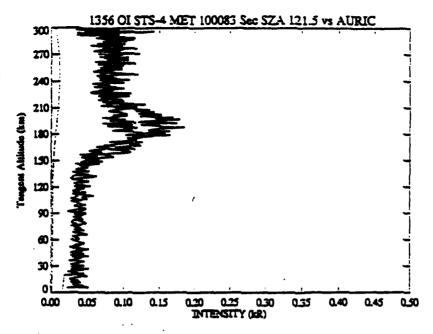


Fig. 9. HUP horizon and scans (up downward) of 1356 A Oxygen emission observed from Space Shuttle STS-4 SZA of 121.5 deg. and overlaid with AURIC prediction (broken line).

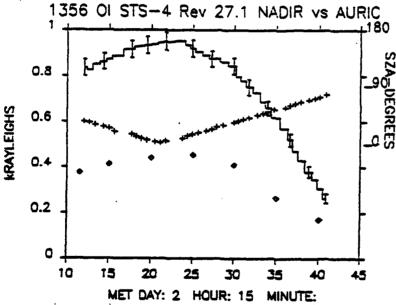
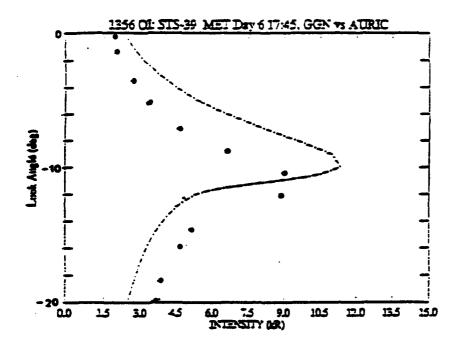


Fig. 10. Peak intensities of the 1356 A oxygen line observed from Space Shuttle STS-4 looking the nadir direction overlaid with SZA and AURIC predictions.

between the observed and the predicted are consistant with that observed in Figures 5 to 8 for the lower tangent altitudes.

Figures 11 and 12 shows the peak 1356 Å line intensities (circles) for the HUP instrument scan angles uncorrected for shuttle attitude while scanning in angle and wavelength on STS-39. The orbiter was in a nose down attitude. We are proceeding to analyze the data for true look angle and tangent height. Here the predicted and observed intensities are in agreement at about 11 kr. The increase in overall brightness compared with data from the earlier flight taken at similar solar zenith angles reflects the fact that these data were obtained at a more active part of the solar cycle, as indicated by the larger 10.7cm solar index value.



Peak Fig. 11. intensities of the Oxygen 1356 A line emission (SZA = 23) vs uncorrected look angle observed from Space Shuttle STS-39 and overlaid with the prediction AURIC (broken line).

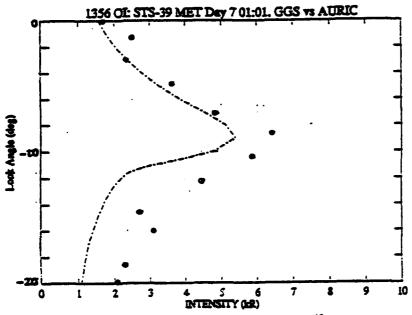


Fig. 12. Peak intensities of the Oxygen 1356 A line emission (SZA = 61) vs uncorrected look angle observed from Space Shuttle STS-39 and overlaid with the AURIC prediction (broken line).

The deduced intensity of the total LBH band system as measured by HUP looking toward the earth from STS-39 is shown in Figure 13. That portion of the measured spectrum between 1376 Å and 1474 Å was integrated and used to represent the entire band system. The bands in that wavelength interval contain 18.6% of the total band strength of the system so the correction was straightforward. Also shown on the figure is an AURIC prediction of the total LBH intensity (small circular point at hour 2, minute 45) for the conditions at the time of the measurement.

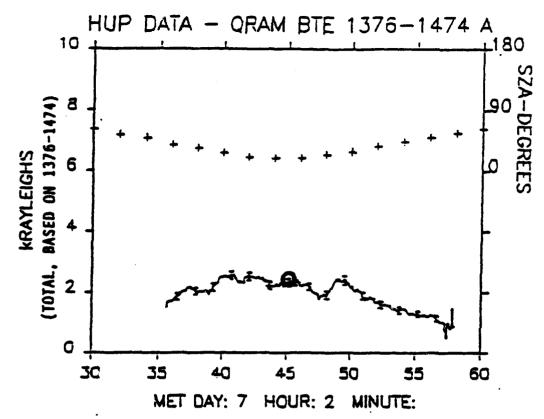


Figure 13. Radiance measured by HUP looking down from space shuttle STS-39 for LBH bands summed between 1376 Å and 1474 Å and corrected for the total LBH system. The AURIC prediction for one point is indicated by a small circle.

## 5. CONCLUSION

Even from this brief comparison it is evident that for well behaved daytime conditions in the ionosphere current state of the art ultraviolet radiance models are capable of providing excellent specification of atmospheric radiance. Future work will concentrate on extending the range of conditions for which experimental data can provide validation for the models.